

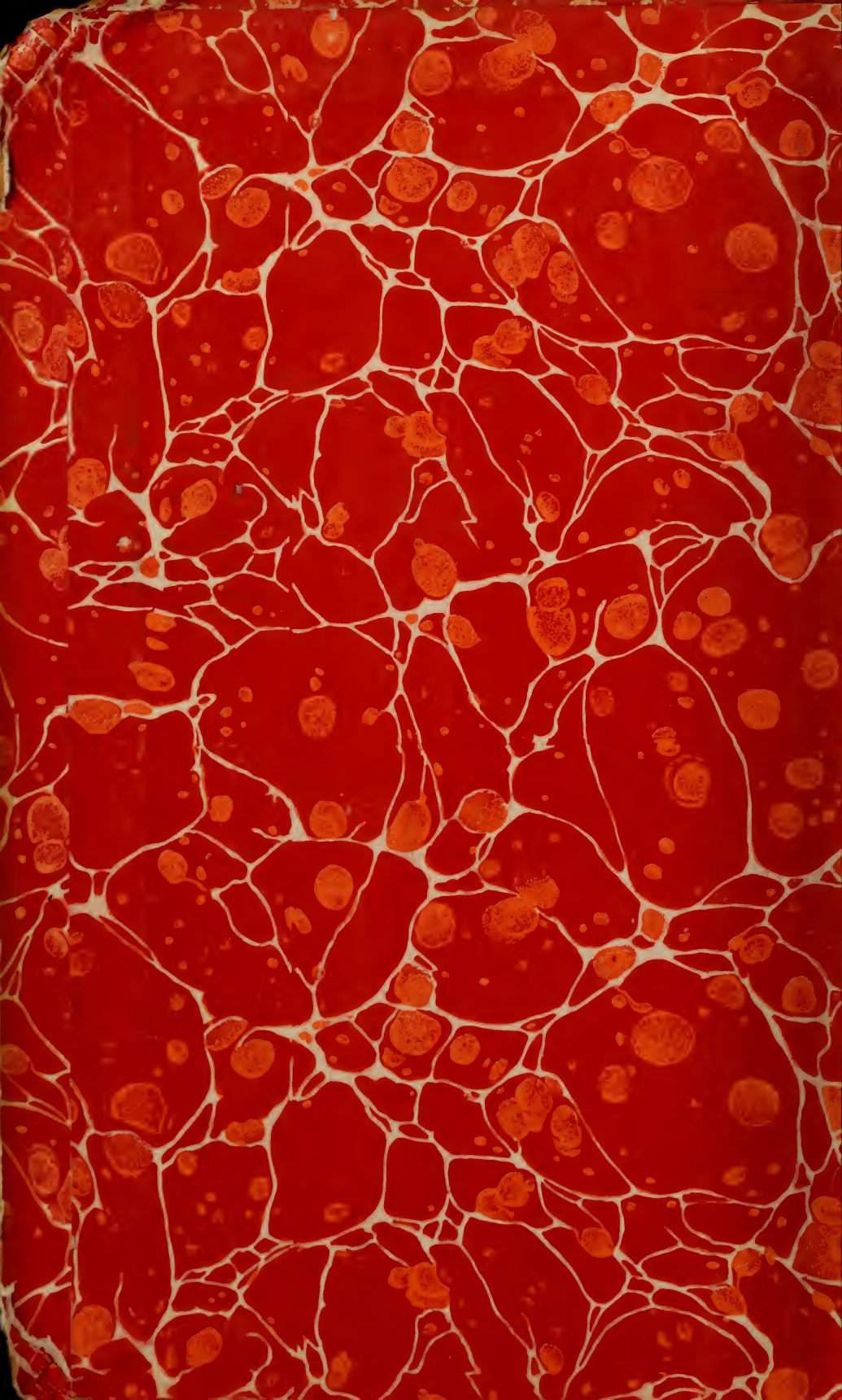
A11102 128969

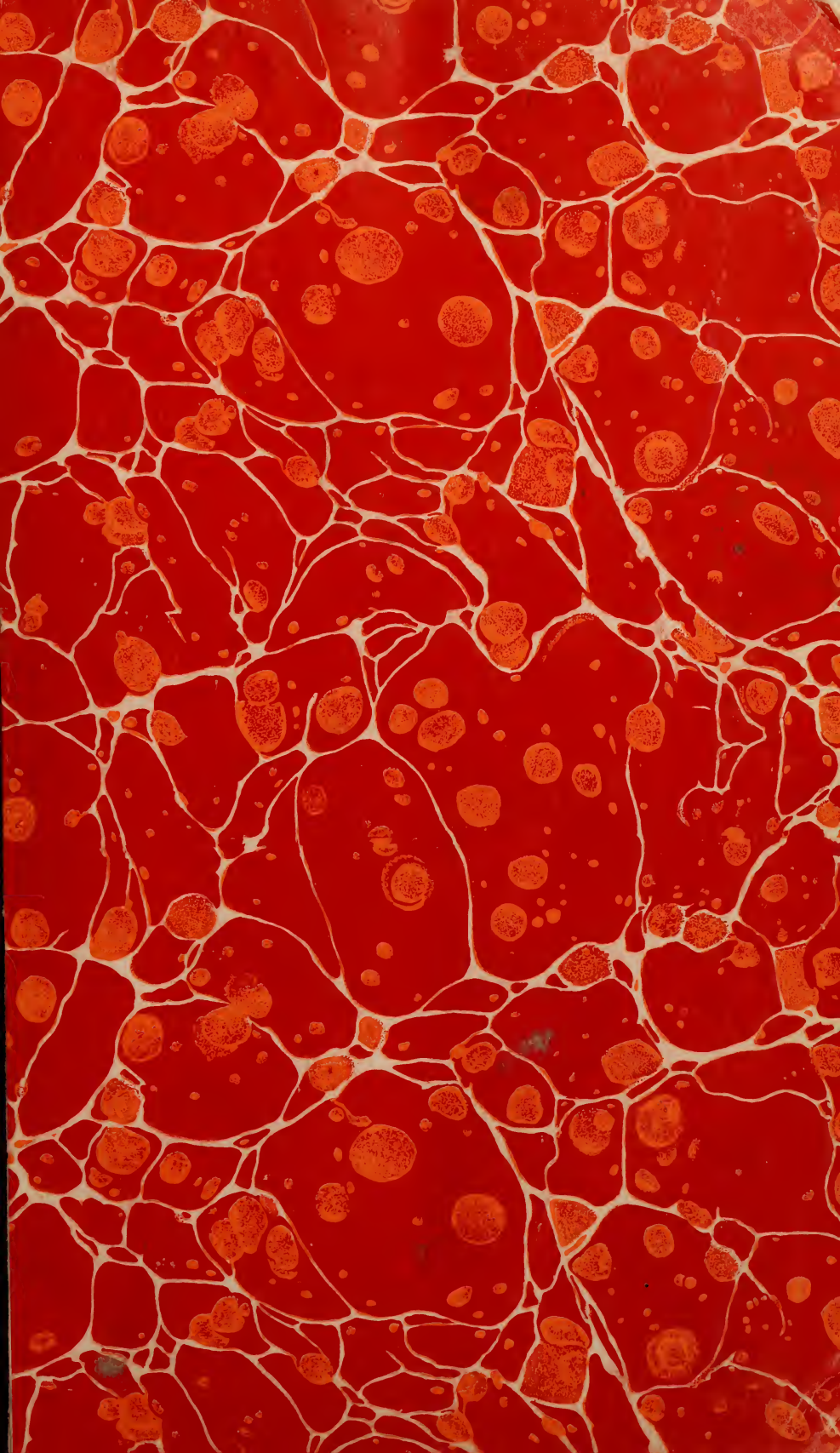
NAT'L INST OF STANDARDS & TECH R.I.C.



A11102128969

/Bureau of Standards Journal of research
QC1 .U52 V10:1933 C.1 NBS-PUB-C 1928









FURTHER DESCRIPTION AND ANALYSIS OF THE FIRST SPECTRUM OF XENON

By C. J. Humphreys and W. F. Meggers

ABSTRACT

An entirely new description of the first spectrum of Xenon has been completed. Geissler tubes with a small bore capillary, exposed end on, were used as sources and afforded sufficient intensity to permit the use of spectrographs of high resolving power to a greater extent than in the preliminary investigation previously reported. The new description covers the region between 3,340 and 11,140 Å. The spectrum has been photographed by aid of a Rowland 20,000-line per inch grating over the range included between 3,440 and 9,800 Å. Interferometer measurements have been made on 130 lines, including 45 reported in a previous paper. Several new lines have been found in the ultra-violet between 3,340 and 3,440 where the use of a source of higher intensity has led to the detection of several higher members of the $1s_5$ - mp_i series. Most of the new data have been obtained in the infra-red region. Each of the new types of Eastman red and infra-red sensitive plates, F, N, P, and Q have been exposed to the radiation in the spectral region for which their respective sensitivity is greatest. It is believed that the spectrum has been completely recorded throughout the region investigated.

Of the 538 lines included in the description all but 13 have been classified. Almost all terms predicted by the Hund theory have been obtained. The new term table embodies a small change in the absolute term values indicated by new measurements of greater precision, several newly discovered terms, and extensions to practically all the series. Some rearrangement of the d -terms has been made necessary, due to the elimination of uncertainties in the assignment of j -values. A new f -type sequence, lying very close to mZ , and designated by mT , has been found.

CONTENTS

	Page
I. Introduction.....	139
II. Experimental details.....	140
III. Wave-length measurements.....	141
IV. Discussion of Xe I terms.....	146

I. INTRODUCTION

A preliminary description and analysis of the first spectrum of xenon was published in 1929.¹ As in the case of the investigation of Kr I which was undertaken simultaneously, it was expected to give in later publications extensions of the series to higher members permitted by the detection of fainter lines, interferometer comparisons of wave lengths, and measurements of hyperfine structures. Interferometer measurements of 45 xenon lines,² and measurements of the hyperfine structures of 6 xenon lines,³ are presented in two papers by

¹ W. F. Meggers, T. L. deBruin, and C. J. Humphreys, B. S. Jour. Research, vol. 3 (RP115), p. 731, 1929.

² C. J. Humphreys, B. S. Jour. Research, vol. 5 (RP245), p. 1041, 1930.

³ C. J. Humphreys, B. S. Jour. Research, vol. 7 (RP351), p. 453, 1931.

one of us in which the results of a similar investigation of Kr I are given. In the case of the first spectrum of krypton it was found that a large number of new lines could be recorded by using as a source a Geissler tube exposed end on. A much more intense illumination of the slit could be obtained in this manner than with the usual type of tube, the construction of which permits only side-on exposures. Extended and improved measurements of Kr I are contained in a paper published about a year ago.⁴ The present paper represents a more complete description of Xe I than could then be made in the case of krypton because in the meantime new photosensitizing dyes have been discovered which have greatly extended both the speed of recording and the upper limit of the range of investigation in the infra-red region. All the rare gas spectra have been reinvestigated, using these specially sensitized plates, and "end on" Geissler tubes (from Robert Goetze in Leipzig). The results for those other than xenon are being published in separate papers. The use of more intense sources and more sensitive plates has made possible additional interferometer comparisons, the results of which for Xe I lines are incorporated in the present description. The application of these superior infra-red sensitive photographic plates to spectrographic investigations has created a demand for accurate standards of wave length in the extended range now accessible and it appears that the emission spectra of the noble gases will be useful for this purpose. The results for Xe I lines presented here may be regarded as a contribution to standard wave lengths, since many of the lines have been measured interferometrically either in terms of neon standards or of sharp xenon lines so determined and enough of the spectral terms have thus been accurately fixed in relative value to guarantee the calculated values of the remaining classified lines to within 0.01 Å.

Since the publication of our first paper, the spectrum of Xe I has been investigated by Gremmer⁵ and by Rasmussen.⁶ Gremmer's classification is essentially the same as ours. He succeeded in obtaining 16 classified lines not recorded in our first description in the ultra-violet region. These are in all cases higher members of known series. Rasmussen found a considerable number of lines in the infra-red, gave classifications for 80 additional lines, and introduced a number of changes in term designations which are confirmed by the present investigation.

II. EXPERIMENTAL DETAILS

New spectograms have been made over the entire region of the Xe I spectrum accessible to photography in air. The Hilger E1 spectrograph was used to describe the ultra-violet region. The last recorded line appears at 3,340 Å. There is no Xe I radiation to be expected between the $1s_5$ series limit at 3,250 Å and the Schumann region. Between 3,440 and 9,800 Å the Rowland grating with 20,000 lines per inch was used, which gives a scale of 3.7 Å/mm in the first order spectrum. The description previously published (RP115) was made for the most part with a grating of 7,500 lines per inch, giving a scale of 10.4 Å/mm. The region between 6,600 and 8,600 Å was also photographed with a scale of 2.7 Å/mm, given by a plane grating with 20,000 lines per inch and lenses of 16 feet focal length. The published obser-

⁴ W. F. Meggers, T. L. deBruin, C. J. Humphreys, B. S. Jour. Research, vol. 7 (RP 364), p. 643, 1931.

⁵ W. Gremmer, Zeit. f. Physik, vol. 59, p. 154, 1930.

⁶ E. Rasmussen, Zeit. f. Physik, vol. 73, p. 779, 1932.

vations beyond 9,800 Å have been made with the 7,500-line-per-inch grating, although the infra-red region was also photographed by aid of a Hilger E2 prism spectograph with glass optics. Interferometer measurements have been made of the wave lengths of 130 Xe lines including those reported in RP245. For lines shown by the interference observations to have hyperfine structure, the wave length of the strongest component only is here reported. A pair of crystal quartz interferometer plates of 60 mm aperture, silvered by the evaporation method, have been used in the more recent work. The fixed étalon method has been employed, using invar separators of 3, 6, 10, 15, and 25 mm length.

In the red and infra-red regions, plates prepared with the new types of sensitizers developed by the Eastman laboratories and designated F, N, P, and Q have been used. The characteristics of these special plates are discussed in a recent paper by Mees.⁷ We have used the F plate to cover the region from 5,000 to 6,600 Å, the N plate from 6,500 to 8,500, the P plate between 8,000 and 9,000, and the Q plate above 9,000. There is considerable overlapping of the regions of sensitivity of these emulsions except that the F type falls off rather abruptly above 6,800 Å. The N plate shows two maxima near 7,000 and 8,200 Å. The respective maxima of the P and Q types lie at 8,600 and 9,700 Å, respectively. Exposures of from 2 to 24 hours have been made. One exposure was prolonged to 40 hours in order to observe the resolution of the pair of lines at 9,374 Å with the Rowland grating, which was necessary to establish experimentally the existence of the *mT* series as discussed below. The stronger lines are considerably overexposed in the case of the longer exposures and, in the work with the Rowland grating, show from 1 to 8 orders of Rowland ghosts. It is believed that the spectrum has been completely recorded in the region studied. Radiometric investigation of still longer waves is highly desirable.

III. WAVE-LENGTH MEASUREMENTS

The wave-length measurements have been made relative to the international standards in the iron arc spectrum. The results from each spectrogram were brought to the scale of the interference measurements by applying a correction where a systematic difference was noted. While such corrections were usually of the order of from 0.01 to 0.02 Å, it was found necessary to apply a correction as great as 0.1 Å in the region above 9,000 Å. The cause of this displacement between the xenon, and second order comparison spectrum of iron, too great to be accounted for by atmospheric dispersion, appears to lie in some unexplained temperature displacement of the plate or mounting during prolonged exposure. Wave lengths which have been determined by interference methods are given to thousandths or ten-thousandths of an angstrom unit and can thus be distinguished from those lines for which only prism or grating measurements are available. The vacuum wave numbers corresponding to the wave lengths determined by interference methods have been computed directly from the atmospheric dispersion formula of Meggers and Peters.⁸ The use of spectrographs of higher dispersion and resolving power, together with

⁷ C. E. K. Mees, J. Opt. Soc. Am., vol. 22, p. 204, 1932.

⁸ W. F. Meggers and C. G. Peters, B. S. Sci. papers, vol. 14 (S. P. 327), p. 722; 1918.

the application of interference methods to one-fourth of the lines, has not only greatly increased the accuracy of the wave-length determinations, but has made possible the resolution of nearly all close pairs of lines. Such pairs occur at 6,111, 6,163, 6,206, 6,840, 7,584, 7,642, 7,664, 8,348, 8,952, and 9,374 Å. In a few cases, notably 7,119 Å, lines formerly regarded as double for purposes of classification are now shown to be single.

The list of Xe I wave lengths as finally adopted is presented in Table 1. After eliminating known lines due to Xe II, faint unclassified lines of doubtful origin, and identifiable Rowland ghosts, we have retained 538 lines. The real lines were in practically all cases distinguished from false lines or ghosts by repeated observations with different gratings. With but 13 exceptions, all of the lines are accounted for by combinations of identified spectral terms. The calculated values of the wave numbers of all classified lines, computed from the experimental terms, are also given in Table 1, and indicate the degree of accuracy of the wave-length determinations.

TABLE 1.—List of Xe I lines

Intensity	Wave length	Wave number observed	Combination	Wave number calculated	Intensity	Wave length	Wave number observed	Combination	Wave number calculated
1—	3,340.04	29,931.18	1s ₅ —11p ₃	1.2	1	3,742.22	26,714.54	1s ₄ —6V	.752
1—	3,348.63	29,854.41	1s ₅ —11U	3.97	10	3,745.38	26,692.00	1s ₄ —6Y	1.95
1h—	3,358.17	29,769.60	1s ₅ —10p ₃	.60	4	3,745.69	26,689.79	1s ₄ —6X	.750
1h—	3,358.96	29,762.60	1s ₅ —10p ₂	.60	3	3,795.95	26,336.42	1s ₅ —5V	.437
1h	3,370.34	29,662.11	1s ₅ —10U	1.95	40	3,796.30	26,333.99	1s ₅ —5U	.95
1	3,383.20	29,549.36	1s ₅ —9p ₃	.36	30	3,801.39	26,298.73	1s ₅ —5Y	.717
1	3,384.36	29,539.23	1s ₅ —9p ₂	.23	3	3,801.90	26,295.20	1s ₅ —5X	.152
2	3,400.07	29,402.75	1s ₅ —9U	.19	30	3,809.84	26,240.40	1s ₄ —5p ₃	.35
1—	3,400.79	29,396.53	1s ₅ —9Y	5.73	10	3,823.74	26,145.02	1s ₄ —5p ₂	.09
2	3,418.37	29,245.35	1s ₅ —8p ₃	.35	15	3,826.86	26,123.70	1s ₄ —5p ₁	.70
2	3,420.00	29,231.41	1s ₅ —8p ₂	.41	2	3,835.6	26,064.2	1s ₄ —5p ₀	5.00
3	3,442.66	29,039.01	1s ₅ —8U	8.77	2	3,942.29	25,358.81	1s ₄ —5V	.821
1	3,443.83	29,029.15	1s ₅ —8Y	8.77	60	3,948.163	25,321.095	1s ₄ —5Y	.101
4	3,469.81	28,811.80	1s ₅ —7p ₃	.86	10	3,948.72	25,317.52	1s ₄ —5X	.536
4	3,472.36	28,790.64	1s ₅ —7p ₂	.64	120	3,950.925	25,303.393	1s ₄ —4p ₃	.395
1	3,496.86	28,688.93	1s ₄ —9p ₃	.93	6	3,956.85	25,265.50	1s ₄ —4p ₁	.541
5	3,506.74	28,508.39	1s ₄ —7U	.43	200	3,967.541	25,197.423	1s ₄ —4p ₂	.423
2	3,508.42	28,494.74	1s ₄ —7Y	5.032	40	3,974.417	25,153.835	1s ₄ —4p ₀	.837
2	3,517.90	28,417.95	1s ₄ —9Y	8.12	30	3,985.202	25,085.762	1s ₄ —4p ₁₀	.762
2	3,533.48	28,292.66	1s ₄ —8p ₃	.66	100	4,078.8207	24,509.996	1s ₄ —4p ₅	.996
1—	3,536.61	28,267.62	1s ₄ —8p ₂	.74	60	4,109.7093	24,325.782	1s ₄ —4p ₆	.779
1	3,537.35	28,261.70	1s ₄ —8p ₁	.70	80	4,116.1151	24,287.925	1s ₄ —4p ₇	.925
10	3,549.86	28,162.11	1s ₄ —6p ₃	.05	20	4,135.1337	24,176.222	1s ₄ —4p ₉	.221
10	3,554.04	28,128.99	1s ₄ —6p ₂	.99	2	4,146.78	24,108.32	1s ₄ —4p ₁₀	.146
1	3,555.92	28,114.12	1s ₅ —6p ₃	.12	20	4,193.01	23,842.52	1s ₅ —4V	.536
3	3,563.80	28,051.95	1s ₄ —8Y	2.16	150	4,193.5296	23,839.567	1s ₅ —4U	.567
4	3,587.02	27,870.37	1s ₄ —7p ₃	.37	50	4,203.6945	23,781.922	1s ₄ —4Y	.921
1	3,591.67	27,834.29	1s ₄ —7p ₂	.25	10	4,205.404	23,772.254	1s ₄ —4X	.253
2	3,592.80	27,825.53	1s ₄ —7p ₁	.53	20	4,372.287	22,864.925	1s ₄ —4V	.920
15	3,610.32	27,690.51	1s ₅ —6U	.53	100	4,383.9092	22,804.306	1s ₄ —4Y	.305
8	3,613.06	27,669.50	1s ₅ —6Y	.56	70	4,385.7693	22,794.635	1s ₄ —4X	.637
6	3,633.06	27,517.19	1s ₄ —7Y	4.16	500	4,500.9772	22,211.188	1s ₅ —2p ₂	.186
10	3,669.91	27,240.90	1s ₄ —6p ₂	.90	400	4,524.6805	22,094.833	1s ₅ —2p ₃	.833
2	3,677.54	27,184.38	1s ₄ —6p ₁	.44	2	4,576.60	21,844.18		
4	3,679.31	27,171.30	1s ₄ —6p ₀	.30	300	4,582.7474	21,814.878	1s ₄ —2p ₁	.875
40	3,685.90	27,122.72	1s ₅ —5p ₃	.70	100	4,611.8896	21,677.034	1s ₅ —3p ₇	.034
1	3,688.80	27,101.40	1s ₅ —5p ₁	.32	1,000	4,624.2757	21,618.973	1s ₅ —3p ₆	.973
40	3,693.49	27,066.99	1s ₅ —5p ₂	.99	2,000	4,671.226	21,401.685	1s ₅ —3p ₈	.685
4	3,696.82	27,042.61	1s ₅ —5p ₀	.61	100	4,690.9711	21,311.603	1s ₅ —2p ₄	.600
2	3,702.74	26,999.37	1s ₅ —5p ₁₀	.41	300	4,697.020	21,284.155	1s ₅ —3p ₉	.154

TABLE 1.—List of Xe I lines—Continued

Inten- sity	Wave length	Wave number observed	Combina- tion	Wave number calcu- lated	Inten- sity	Wave length	Wave number observed	Combina- tion	Wave number calcu- lated
5	4,708.21	21,233.57	1s ₁ -2p ₂	.570	3	5,607.99	17,826.76	2p ₈ -8d ₁ '	.820
600	4,734.1524	21,117.217	1s ₁ -2p ₃	.217	15	5,612.65	17,811.96	2p ₈ -8d ₁ ''	2.01
150	4,792.6192	20,859.605	1s ₁ -3p ₁₀	.605	80	5,618.878	17,792.220	2p ₈ -8d ₄	.220
500	4,807.019	20,797.118	1s ₁ -3p ₃	.118	1	5,621.24	17,784.74	2p ₈ -8d ₃	.805
400	4,829.709	20,699.415	1s ₁ -3p ₇	.418	5	5,646.19	17,706.16	2p ₈ -7s ₃	.17
300	4,843.294	20,641.357	1s ₁ -3p ₆	.357	2h	5,652.84	17,685.33	2p ₇ -10d ₁ ''	.31
500	4,916.508	20,333.981	1s ₁ -2p ₄	.984	1h—	5,654.31	17,680.73	2p ₈ -8d ₆	.77
500	4,923.1522	20,306.538	1s ₁ -3p ₉	.538	5	5,664.46	17,649.05	2p ₆ -11d ₁ '	.05
3h	5,023.88	19,899.40	1s ₃ -8X	.47	40	5,688.373	17,574.855	1s ₂ -6V	.855
200	5,028.2796	19,881.989	1s ₃ -3p ₁₀	.989	100	5,695.750	17,552.094	1s ₂ -6Y	.094
10	5,162.711	19,364.290	1s ₃ -7X	.290	80	5,696.479	17,549.845	1s ₂ -6X	.851
1h	5,164.39	19,357.99	2p ₁₀ -10d ₃	.94	8	5,698.54	17,543.50	2p ₈ -8d ₁ '	.561
1h	5,167.30	19,347.09	2p ₁₀ -10d ₅	.08	1	5,703.34	17,528.73	2p ₈ -8d ₁ ''	.75
2h	5,185.85	19,277.89	1s ₂ -9Y	8.22	3	5,706.87	17,517.89	2p ₁₀ -5s ₄	.954
1	5,206.07	19,203.01	2p ₁₀ -8s ₃	2.99	10h	5,709.80	17,508.90	2p ₈ -8d ₄	.961
4h	5,245.27	19,059.51	2p ₁₀ -9d ₃	.48	2	5,712.21	17,501.52	2p ₈ -8d ₃	.546
4h	5,248.98	19,046.03	2p ₁₀ -9d ₅	.02	70	5,715.716	17,490.780	2p ₁₀ -5s ₃	.783
2h	5,251.89	19,035.48	1s ₃ -6p ₁₀	.59	80	5,716.252	17,489.140	2p ₈ -8d ₄ '	.140
1h	5,273.48	18,957.55	1s ₂ -8V	.31	15h	5,722.14	17,471.15	2p ₆ -6s ₄	.18
2h	5,283.30	18,922.31	1s ₂ -8Y	.26	1h	5,723.26	17,467.72	2p ₇ -9d ₂	.74
4h	5,286.11	18,912.25	1s ₂ -8X	.20	4	5,726.10	17,459.06	2p ₆ -6s ₃	.131
3h	5,286.38	18,911.29	2p ₁₀ -7s ₃	.08	4h	5,733.48	17,436.69	2p ₆ -10d ₁ '	.60
3	5,306.37	18,840.05	2p ₈ -11d ₁ ''	.78	6	5,740.17	17,416.27	2p ₁₀ -6d ₂	.29
1h—	5,335.91	18,735.73	2p ₈ -11d ₄ '	.80	1h	5,740.73	17,414.57	2p ₆ -10d ₃	.62
2h	5,337.89	18,728.80	2p ₈ -13d ₄ '	.66	5h	5,748.20	17,391.94	2p ₇ -9d ₁ '	2.00
1	5,356.50	18,662.68	2p ₁₀ -8d ₁ ''	.739	1h	5,754.60	17,372.59	2p ₇ -9d ₃	.59
15	5,362.244	18,643.739	2p ₁₀ -8d ₃	.459	1h	5,792.26	17,259.64	2p ₈ -8s ₃	.67
30	5,364.626	18,635.459	2p ₁₀ -8d ₅	.85	15	5,807.311	17,214.912	2p ₈ -7d ₁ '	.922
6	5,367.03	18,627.11	2p ₈ -12d ₄ '	.119	60	5,814.505	17,193.614	2p ₈ -7d ₁ ''	.616
1h—	5,373.74	18,603.85	1s ₃ -6X	.443	25	5,820.52	17,175.84	2p ₈ -6s ₃	.907
100	5,392.795	18,538.119	2p ₁₀ -8d ₆	.55	300	5,823.890	17,165.906	1s ₃ -5X	.907
20	5,394.738	18,531.443	2p ₈ -10d ₁ ''	.54	150	5,824.800	17,163.225	2p ₆ -7d ₄	.225
1h	5,397.63	18,521.51	2p ₈ -10s ₃	.84	1	5,827.72	17,154.63	2p ₈ -7d ₃	.636
4h	5,400.45	18,511.84	2p ₈ -10d ₄ '	.08	20h	5,830.63	17,146.06	2p ₆ -9d ₁	.10
5	5,418.02	18,451.81	2p ₈ -8s ₃	.27	4h	5,840.83	17,116.12	2p ₈ -9d ₃	.16
2h	5,421.76	18,439.08	2p ₈ -11d ₁ '	.13	5	5,843.43	17,108.51	2p ₈ -7d ₃	.610
5h	5,435.60	18,392.13	1s ₂ -7V	.519	1	5,845.46	17,102.56	2p ₆ -9d ₃	.70
30	5,439.923	18,377.519	1s ₂ -7Y	6.022	2	5,846.21	17,100.37	1s ₂ -5p ₃	.45
15	5,440.39	18,375.94	1s ₂ -7X	.82	3h	5,849.85	17,089.73	2p ₇ -8d ₂	.74
1h	5,444.87	18,360.82	2p ₈ -8s ₄	.24	15	5,856.309	17,070.300	2p ₁₀ -6d ₁ '	.300
1h	5,454.54	18,328.27	2p ₈ -9s ₃	.83	100	5,875.018	17,016.519	2p ₁₀ -6d ₃	.519
2	5,456.45	18,321.85	2p ₁₀ -6s ₄	.820	6	5,878.92	17,005.23	1s ₂ -5p ₆	.19
15	5,460.037	18,309.820	2p ₁₀ -6s ₃	.01	20	5,889.12	16,975.77	2p ₇ -8d ₁ '	.77
1h	5,479.12	18,246.05	2p ₈ -10d ₁ '	.77	100	5,894.958	16,958.874	2p ₁₀ -6d ₃	.874
1h	5,481.33	18,238.69	2p ₈ -9d ₁ '	8.60	2h	5,895.62	16,957.05	2p ₇ -8d ₃	6.850
1h	5,484.16	18,229.28	2p ₈ -10d ₄	.24	8	5,898.56	16,948.60	2p ₇ -8d ₃	.570
4h	5,484.46	18,228.28	2p ₈ -9d ₁ ''	.75	20	5,904.462	16,931.663	2p ₈ -7d ₁ '	.663
6h	5,487.03	18,219.75	2p ₈ -10d ₄ '	.684	3	5,906.76	16,925.07	1s ₂ -5p ₉	.10
20h	5,488.555	18,214.684	2p ₈ -9d ₄	.08	5	5,911.90	16,910.36	2p ₈ -7d ₁ ''	.357
2h	5,532.78	18,069.10	2p ₈ -8s ₅	.34	4	5,916.65	16,896.78	2p ₆ -7s ₃	.76
3h	5,540.38	18,044.30	1s ₂ -6p ₆	.270	10	5,921.85	16,881.95	1s ₂ -5p ₁₀	.90
50	5,552.385	18,005.291	2p ₁₀ -7d ₃	.291	20	5,922.550	16,879.953	2p ₈ -7d ₄	.966
3h	5,553.10	18,002.97	2p ₈ -7s ₄	.97	6	5,925.56	16,871.38	2p ₈ -7d ₃	.377
1—	5,555.06	17,996.62	1s ₂ -6p ₉	.61	80	5,931.241	16,855.219	2p ₁₀ -6d ₃	.220
2	5,557.28	17,989.43	2p ₈ -7s ₃	.43	100	5,934.172	16,846.894	2p ₈ -7d ₄ '	.894
2	5,563.50	17,969.32	1s ₂ -6p ₁₀	.264	1—	5,970.41	16,744.64	3d ₅ -10V	.64
5h	5,566.22	17,960.54	2p ₁₀ -7d ₅	.51	1h	5,972.82	16,737.88	3d ₅ -10X, Y	.153
100	5,566.615	17,959.264	2p ₈ -9d ₁ '	.425	40	5,974.152	16,734.153	2p ₈ -8d ₁ '	.57
2h	5,567.77	17,955.54	2p ₈ -9d ₁ '	.50	2	5,978.29	16,722.57	3d ₄ '-11T, Z	.34
2h	5,575.27	17,931.38	2p ₈ -9d ₄	.50	1	5,979.42	16,719.41	2p ₆ -8d ₁ '	.418
40	5,579.28	17,918.50	2p ₈ -8d ₃	.458	4	5,986.23	16,700.39	2p ₆ -8d ₃	.138
50	5,581.784	17,910.458	2p ₁₀ -7d ₅	.54	20	5,989.18	16,692.16	2p ₆ -8d ₂	.300
1—	5,585.18	17,899.57	2p ₇ -11d ₁ ''	.17	30	5,998.115	16,687.300	2p ₈ -5s ₃	.128
6	5,594.37	17,870.16	1s ₂ -5p ₁₀	.425	15	6,007.909	16,640.128	2p ₈ -5s ₅	.94
					8	6,009.78	16,634.94	2p ₇ -6s ₄	.94

TABLE 1.—List of Xe I lines—Continued

Inten- sity	Wave length	Wave number observed	Combina- tion	Wave num- ber calcu- lated	Inten- sity	Wave length	Wave number observed	Combina- tion	Wave num- ber calcu- lated
1h	6,014.10	16,623.00	2p ₇ -6s ₅	2.931	20	6,430.155	15,547.436	2p ₆ -5s ₅	.461
1h	6,022.89	16,598.74	3d ₅ -11U	.74	2h	6,448.70	15,502.72	3d ₄ -9W	.71
4	6,026.76	16,588.03	2p ₆ -8d ₆	.10	7	6,450.48	15,498.45	2p ₆ -2s ₂	.45
1h	6,031.36	16,575.43	2p ₅ -10d ₂	.46	10hl	6,451.79	15,495.30	3d ₄ -9Z	.30
2	6,034.92	16,565.65	2p ₉ -6d ₂	.64	3	6,461.50	15,472.02	2p ₅ -6s ₄	.01
10	6,043.38	16,542.46	2p ₇ -7d ₂	.45	300	6,469.705	15,452.394	2p ₁₀ -5d ₃	.394
6h	6,048.00	16,529.82	3d ₄ '-10T, Z	.82	150	6,472.841	15,444.906	2p ₁₀ -5d ₂	.906
1h	6,064.91	16,483.74	3d ₅ -9V	.74	120	6,487.765	15,409.378	2p ₁₀ -5d ₁ '	.383
2h	6,067.52	16,476.65	3d ₅ -9Y	.62	30hl	6,497.43	15,386.46	3d ₄ '-7W	.48
1h	6,067.77	16,475.97	3d ₅ -9X	.98	100	6,498.718	15,383.407	2p ₇ -6d ₁ '	.411
3	6,093.38	16,406.72	3d ₃ -10U	.72	15	6,500.37	15,379.50	2p ₅ -7d ₂	.52
1-	6,095.15	16,401.96	3d ₃ -10X, Y		200h	6,504.18	15,370.49	3d ₄ '-7T, Z	.49
3	6,103.88	16,378.50	2p ₆ -6s ₄	.51	3	6,507.50	15,362.65	1s ₂ -4p ₅	.10
8	6,103.37	16,366.46	2p ₆ -6s ₅	.519	40	6,521.508	15,329.648	2p ₇ -6d ₃	.630
30	6,111.759	16,357.384	2p ₇ -7d ₁ '	.351	100	6,533.159	15,302.312	2p ₆ -4s ₄	.312
40	6,111.951	16,356.870	2p ₆ -5s ₅	.869	40	6,543.360	15,278.456	2p ₆ -4s ₅	.455
10	6,114.86	16,349.09	2p ₁₀ -2s ₂	.10	20	6,546.12	15,272.01	2p ₇ -6d ₃	1.985
5	6,123.91	16,324.93	3d ₅ -8X	.96	40	6,553.66	15,254.45	3d ₅ -7V	.43
15	6,126.36	16,318.40	2p ₇ -7d ₃	.401	50hl	6,554.196	15,253.196	3d ₅ -7U	.20
1	6,131.47	16,304.80	2p ₅ -9d ₂	.80	25	6,559.97	15,239.77	3d ₅ -7Y	.80
1h	6,142.13	16,276.50	3d ₄ '-9W	.53	4h	6,560.65	15,238.19	3d ₅ -7X	.30
4	6,143.70	16,272.34	2p ₇ -7d ₅	.375	20	6,583.27	15,185.83	1s ₂ -4p ₆	.882
20hw	6,144.97	16,268.98	3d ₄ '-9T, Z	.98	8	6,590.86	15,168.35	2p ₇ -6d ₅	.331
20	6,152.069	16,250.206	2p ₆ -6d ₁ '	.201	100	6,595.561	15,157.534	2p ₆ -6d ₁ '	.534
3	6,162.16	16,223.60	2p ₇ -7d ₅	.569	4h	6,602.87	15,140.75	3d ₄ -8W	.75
90	6,163.660	16,219.648	2p ₆ -6d ₁ '	.646	30h	6,607.41	15,130.35	3d ₄ -8Z	.30
80	6,163.935	16,218.924	1s ₂ -5V	.924	10	6,608.87	15,127.01	2p ₆ -6d ₁ '	6.979
150	6,178.302	16,181.209	1s ₂ -5Y	.204	12	6,630.44	15,077.80	2p ₆ -6d ₄	.765
120	6,179.665	16,177.640	1s ₂ -5X	.639	50	6,632.464	15,073.198	2p ₆ -6d ₃	.198
300	6,182.420	16,170.431	2p ₆ -6d ₄	.432	3	6,648.75	15,036.28	1s ₂ -4p ₅	.324
3	6,184.16	16,165.88	2p ₆ -6d ₃	.865	20	6,657.92	15,015.57	2p ₆ -6d ₅	.553
20	6,189.10	16,152.98	2p ₁₀ -4s ₄	.966	4	6,664.85	14,999.95	2p ₅ -4s ₅	.196
4h	6,191.40	16,146.98	3d ₅ -9U	.96	60	6,666.965	14,995.196	2p ₁₀ -5d ₅	.800
1h	6,193.89	16,140.49	3d ₅ -9Y	.50	150	6,668.920	14,990.800	1s ₂ -4p ₁₀	.249
1	6,195.49	16,136.32	1s ₂ -4p ₇	.296	25	6,678.972	14,968.239	3d ₅ -6X	.613
100	6,198.260	16,129.108	2p ₁₀ -4s ₅	.109	20	6,681.036	14,906.89	3d ₅ -6p ₅	.82
60	6,200.890	16,122.266	2p ₆ -7d ₁ '	.255	1	6,706.46	14,859.146	2p ₁₀ -5d ₅	.148
3h	6,201.49	16,120.71	3d ₅ -8V	.71	200	6,728.008	14,859.146	3d ₄ -6p ₅	8.89
6h	6,205.35	16,110.68	3d ₅ -8Y	.66	10	6,767.12	14,773.27	3d ₅ -6V	.26
4	6,205.75	16,109.64	3d ₅ -8X	.60	50	6,777.57	14,750.48	3d ₅ -6Y	.49
20	6,206.297	16,108.220	2p ₆ -6d ₅	.220	40	6,778.60	14,748.25	3d ₅ -6X	.251
3	6,209.11	16,100.92	2p ₆ -7d ₁ '	.949	15	6,815.64	14,688.09	2p ₆ -5s ₄	.129
1	6,220.84	16,070.53	2p ₆ -7d ₄	.558	20	6,818.38	14,682.20	2p ₇ -2s ₁	.21
40	6,224.169	16,061.968	2p ₆ -7d ₃	.969	200	6,827.315	14,643.014	1s ₂ -4X	.008
8	6,242.09	16,015.85	2p ₆ -7d ₅	.943	8	6,840.96	14,613.80	2p ₆ -5d ₁ '	.804
50	6,261.212	15,966.942	2p ₅ -6d ₁ '	.942	20	6,841.50	14,612.65	3d ₄ '-7W	.66
40	6,265.301	15,956.522	1s ₂ -4p ₁₀	.517	1	6,844.27	14,606.74	3d ₄ '-7V	.73
1h	6,268.34	15,948.78	3d ₄ -11Z	.78	2h	6,844.84	14,605.52	3d ₄ '-7U	.50
10	6,273.23	15,936.35	2p ₆ -6d ₁ '	.387	60	6,846.613	14,601.740	2p ₆ -5d ₃	.740
4	6,276.99	15,926.81	2p ₅ -8d ₂	.81	50	6,848.82	14,597.04	3d ₄ '-7Z	.05
5h	6,281.81	15,914.53	3d ₄ '-8W	.57	30	6,850.13	14,594.24	2p ₆ -5d ₂	.252
100	6,286.011	15,903.950	3d ₄ '-8T	.95	40	6,860.19	14,572.84	3d ₄ '-6W	.82
50	6,292.649	15,887.173	2p ₆ -6d ₄	.173	20	6,863.20	14,566.45	2p ₆ -6d ₃	.47
15	6,294.45	15,882.63	2p ₆ -6d ₃	.608	5	6,865.58	14,561.40	3d ₄ '-6U	.42
15	6,314.97	15,831.02	2p ₇ -5s ₄	.065	50	6,868.838	14,558.734	2p ₆ -5d ₁ '	.729
500	6,318.062	15,823.273	2p ₆ -6d ₁ '	.273	100	6,872.107	14,547.571	3d ₄ '-6T	.57
2	6,325.81	15,803.89	2p ₇ -5s ₅	.893	300	6,882.155	14,526.332	2p ₆ -5d ₄	.333
20	6,331.50	15,789.09	3d ₆ -7X	.784	30	6,910.82	14,466.08	2p ₇ -4s ₄	.077
40hl	6,333.97	15,783.53	3d ₆ -8U	.54	8	6,922.22	14,442.25	2p ₇ -4s ₅	.220
shl	6,337.53	15,774.64	3d ₆ -8Y	.54	15	6,924.67	14,437.15	3d ₅ -6V	.14
2h	6,344.98	15,756.14	3d ₄ -10Z	.15	100	6,925.53	14,435.35	3d ₅ -6U	.30
20	6,355.77	15,729.40	2p ₇ -6d ₂	.40	50	6,935.62	14,414.35	3d ₅ -6Y	.37
10	6,412.38	15,590.53	3d ₅ -7V	.55	8	6,936.69	14,412.13	3d ₅ -6X	.13
30	6,418.41	15,575.89	3d ₅ -7Y	.92					
30h	6,418.93	15,574.50	2p ₆ -5s ₄	.633					
			3d ₅ -7X	.42					

TABLE 1.—List of Xe I lines—Continued

Intensity	Wave length	Wave number observed	Combina- tion	Wave number calculated	Intensity	Wave length	Wave number observed	Combina- tion	Wave number calculated
1	6,949.76	14,385.02			10	7,666.61	13,039.99	3d ₃ -5X	.92
100	6,976.182	14,330.542	2p ₈ -5d ₁ '	.545	1	7,670.81	13,032.85	2p ₆ -3s ₁ ''	.86
30	6,982.05	14,318.50	2p ₈ -5d ₃	.481	40	7,740.31	12,915.83	2p ₆ -5d ₅	.825
1	6,991.65	14,298.84	3d ₅ -5p ₅	.85	50	7,783.66	12,843.90	3d ₁ ''-6W	.86
4	7,003.10	14,275.46	2p ₈ -5d ₁ ''	.470	15	7,789.42	12,834.40	3d ₁ ''-6V	.30
30	7,019.02	14,243.08	2p ₈ -5d ₄	.073	1	7,790.53	12,832.57	3d ₁ ''-6U	.46
3	7,034.80	14,211.13			100	7,802.651	12,812.636	2p ₆ -3s ₄	.636
20	7,035.53	14,209.66	2p ₆ -4s ₄	.645	10	7,832.98	12,763.02	2p ₁₀ -4d ₂	.00
30	7,047.37	14,185.79	2p ₆ -4s ₅	.788	15	7,841.23	12,749.60	2p ₈ -3s ₁ ''	.60
1h	7,049.07	14,182.36	3d ₅ -5p ₇	.21	100	7,881.320	12,684.745	2p ₆ -3s ₅	.742
1h	7,049.36	14,181.78	3d ₁ ''-8V	.75	300	7,887.395	12,674.975	1s ₂ -2p ₁	.978
3	7,051.06	14,178.36	2p ₁₀ -3s ₁ ''''	.34	40	7,937.41	12,595.11	2p ₈ -5d ₂	.081
1	7,078.46	14,123.48	3d ₅ -5p ₆	.50	4	7,954.22	12,568.49	3d ₅ -4p ₅	.50
500	7,119.598	14,041.872	2p ₈ -5d ₄	.872	500	7,967.341	12,547.792	1s ₃ -3p ₇	.789
15	7,136.57	14,008.48	2p ₆ -5d ₅	.491	8	7,976.03	12,534.12	2p ₆ -3s ₁ ''''	.100
10	7,172.70	13,937.92	3d ₄ '-5p ₈	.88					
15	7,200.79	13,883.54	2p ₁₀ -3s ₁ ''	.51	10	8,003.26	12,491.48	2p ₇ -3s ₁ ''''	.45
5	7,209.14	13,867.46	3d ₅ -5p ₆	.47	100	8,029.67	12,450.39	3d ₄ -5W	.41
1	7,220.24	13,846.14	3d ₃ -5p ₇	.09	10	8,040.56	12,433.53	3d ₄ -5V	.51
3	7,238.20	13,811.79	3d ₅ -5p ₈	.76	15	8,042.18	12,431.02	3d ₄ -5U	.02
20	7,244.94	13,798.94	3d ₄ -6W	9.00	200	8,057.258	12,407.762	3d ₄ -5Z	.75
2	7,249.92	13,789.46	3d ₄ -6V	.44	150	8,061.340	12,401.480	2p ₈ -3s ₅	.483
5h	7,250.87	13,787.66	{ 3d ₅ -5p ₈	.38	2	8,064.94	12,395.94	3d ₄ -5Y	.79
			{ 3d ₃ -6U	.60	1	8,073.99	12,382.05	3d ₅ -4p ₁₀	.011
60	7,257.94	13,774.22	3d ₄ -6Z	.23	3	8,097.24	12,346.50	3d ₅ -4p ₇	.43
20	7,262.54	13,765.50	2p ₇ -5d ₃	.505	100	8,101.98	12,339.27	3d ₁ '-6W	.26
25	7,266.49	13,758.02	2p ₇ -5d ₂	.017	6	8,107.91	12,330.25	3d ₁ '-6V	9.71
40	7,283.961	13,725.018	1s ₂ -4V	.023	15	8,109.46	12,327.89	3d ₁ '-6U	.86
60	7,285.301	13,722.494	2p ₇ -5d ₁ ''	.494	15	8,118.29	12,314.48	3d ₁ '-6Z	.49
5h	7,307.37	13,681.05	3d ₁ '-8W	.01	2	8,123.29	12,306.90	3d ₁ '-6Y	.93
1h	7,313.01	13,670.50	3d ₁ '-8Z	.56	2	8,165.37	12,243.48		
70	7,316.272	13,664.405	1s ₂ -4Y	.408	100	8,171.02	12,235.01	{ 3d ₅ -4p ₆	4.74
20	7,316.87	13,663.28	2p ₁₀ -3s ₄	.290				{ 2p ₆ -3s ₁ ''''	.02
15	7,319.94	13,657.56	3d ₁ ''-7W	.52	1-	8,182.93	12,217.21	3d ₂ -8V	.40
80	7,312.452	13,654.737	1s ₂ -4X	.740	2	8,196.73	12,196.63	2p ₇ -3s ₁ ''	.62
2	7,323.05	13,651.76	2p ₁ '-7V	.59	700	8,206.341	12,182.354	1s ₃ -2p ₄	.355
					10,000	8,231.6348	12,144.921	1s ₅ -2p ₆	.923
50	7,336.480	13,626.767	2p ₇ -3s ₁ '''	.767	500	8,266.519	12,093.671	1s ₂ -2p ₂	.673
40	7,355.58	13,591.38	3d ₅ -5X	.401	7,000	8,280.1163	12,073.811	1s ₄ -2p ₅	.811
100	7,386.002	13,535.397	2p ₁₀ -3s ₅	.396	15	8,297.71	12,048.21	3d ₅ -4p ₆	.16
150	7,393.793	13,521.140	2p ₆ -5d ₁ '	.137	2	8,323.90	12,010.31	3d ₃ -4p ₇	.31
30	7,400.41	13,509.04	2p ₆ -5d ₃	.077	20	8,324.58	12,009.32	2p ₇ -5d ₅	.301
12	7,404.51	13,501.57	2p ₆ -5d ₂	.585	2,000	8,346.823	11,977.318	1s ₂ -2p ₃	.320
3	7,405.77	13,499.26	2p ₆ -2s ₂	.28	60	8,347.45	11,976.42	2p ₇ -3s ₄	.401
20	7,424.05	13,466.03	2p ₆ -5d ₁ ''	.062	40	8,349.05	11,974.12	2p ₁₀ -4d ₁ ''	.10
20	7,441.94	13,433.66	2p ₆ -5d ₄	.665	3	8,371.38	11,942.18	3d ₅ -4p ₈	.19
25	7,451.00	13,417.33	3d ₅ -5V	.33	5	8,372.79	11,940.17	2p ₆ -3s ₁ ''	.19
40	7,472.01	13,379.60	3d ₅ -5Y	.61	20	8,392.37	11,912.32	2p ₆ -4d ₂	.35
25	7,474.01	13,376.02	3d ₅ -5X	.04	5	8,402.03	11,898.62	3d ₅ -4p ₆	.60
20	7,492.23	13,343.49	2p ₈ -3s ₁ '''	.508	2,000	8,409.190	11,888.489	1s ₅ -2p ₇	.491
20	7,501.13	13,327.66	2p ₆ -3s ₁ ''''	.69	10	8,437.55	11,848.53	2p ₇ -3s ₅	.507
8	7,514.54	13,303.88	2p ₇ -5d ₆	.911	1h	8,450.37	11,830.55	3d ₅ -4p ₁₀	.53
3	7,514.96	13,303.13	2p ₅ -4s ₄	.141	1	8,501.02	11,760.07		
40	7,559.79	13,224.25	3d ₄ '-5W	.23	30	8,522.55	11,730.35	1s ₃ -3p ₁₀	.360
6	7,570.93	13,204.79	3d ₄ '-5U	.84	30	8,530.10	11,719.98	2p ₆ -3s ₄	.969
10	7,584.29	13,181.53	3d ₄ '-5Z	.57	2	8,553.97	11,687.27	3d ₂ -7V	.24
200	7,584.680	13,180.849	3d ₄ '-5T	.85	1-	8,564.7	11,672.63	3d ₂ -7Y	.60
6	7,589.61	13,172.28	2p ₇ -5d ₅	.257	200	8,576.01	11,657.24	1s ₂ -3p ₅	.221
1	7,594.36	13,164.04	3d ₄ '-5p ₃	.06	80	8,624.24	11,592.04	2p ₆ -3s ₅	.075
10	7,600.77	13,152.95	3d ₄ '-7W	.92	250	8,643.54	11,559.47	1s ₂ -3p ₇	.521
2h	7,604.97	13,148.68	3d ₄ '-7U	.80	100	8,692.20	11,501.41	1s ₂ -3p ₆	.460
5	7,608.46	13,139.65	3d ₄ '-5p ₉	.68	200	8,696.86	11,495.25	3d ₁ ''-5W	.27
3	7,609.82	13,137.30	3d ₄ '-7Z	.31	40	8,709.64	11,478.38	3d ₁ ''-5V	.37
500	7,642.025	13,081.940	1s ₃ -2p ₂	.941	2	8,711.54	11,475.88	3d ₁ ''-5U	.88
	7,642.30	13,081.47	3d ₃ -5V	.21	300	8,739.39	11,439.31	2p ₁₀ -4d ₃	.31
100	7,643.91	13,078.71	3d ₃ -5U	.72	100	8,758.20	11,414.74	2p ₆ -4d ₁ '	.75
10	7,664.02	13,044.40	2p ₆ -3s ₁ ''''	.43	5,000	8,819.412	11,335.515	1s ₅ -2p ₈	.515
30	7,664.56	13,043.48	3d ₃ -5Y	.49					

TABLE 1.—List of Xe I lines—Continued

Intensity	Wave length	Wave number observed	Combination	Wave number calculated	Intensity	Wave length	Wave number observed	Combination	Wave number calculated
1	8,851.44	11,294.50	$3d_1-4p_3$.49	3	9,605.80	10,407.53	$3d_1''-4p_7$.47
300	8,862.32	11,280.63	$2p_{10}-4d_5$.63	1	9,614.43	10,398.18		
10	8,885.71	11,250.94	$3d_1-4p_9$.90	1	9,616.95	10,395.46		
200	8,908.73	11,221.87	$2p_{10}-4d_6$.87	1	9,668.94	10,339.56		
200	8,930.83	11,194.10	$1s_2-2p_4$	4.087	150	9,685.32	10,322.08	$2p_6-4d_1'$.08
1,000	8,952.254	11,167.309	$1s_1-2p_3$.307	20	9,700.99	10,305.40	$2p_8-4d_3$.40
50	8,952.78	11,166.65	$1s_2-3p_3$.641	2	9,710.03	10,295.81	$3d_1''-4p_9$.76
100	8,981.05	11,131.50	$2p_8-4d_1'$.49	100	9,718.16	10,287.20	$2p_7-4d_1''$.21
200	8,987.57	11,123.43	$2p_9-4d_1''$.45	2,000	9,799.699	10,201.600	$1s_5-2p_{10}$.602
30	9,025.98	11,076.09	$2p_7-4d_3$.11	3,000	9,923.192	10,074.643	$1s_1-2p_9$.640
50	9,032.18	11,068.49	$3d_5-4X$.502	10	9,966.58	10,030.78	$2p_6-4d_1''$.78
400	9,045.446	11,052.256	$1s_5-2p_2$.256	50	10,023.72	9,973.60	$3d_1-4W$.58
50	9,096.13	10,990.67	$3d_1'-5W$.67	1	10,056.84	9,940.76		
4	9,112.24	10,971.24	$3d_1'-5U$.28	5	10,057.96	9,939.66	$3d_1-4V$.60
3	9,131.59	10,947.99	$3d_1'-5Z$.01	10	10,060.96	9,936.69	$3d_1-4U$.63
2	9,141.8	10,935.8	$3d_1'-5Y$	6.05	20	10,084.79	9,913.21	$2p_8-4d_2$.18
20	9,152.12	10,923.44	$3d_5-4V$.42	80	10,107.34	9,891.10	$3d_1-4Z$.08
2	9,158.38	10,915.97			1	10,119.8	9,878.88	$3d_1-4Y$.99
500	9,162.654	10,910.877	$1s_4-2p_7$.875	20	10,125.47	9,873.38	$3d_5-2p_1$.35
100	9,167.52	10,905.09	$2p_9-4d_1$.09	10	10,188.36	9,812.44	$2p_7-4d_1$.42
2	9,197.18	10,869.92	$3d_5-6V$.96	20	10,251.07	9,752.41	$2p_7-4d_3$.42
30	9,203.20	10,862.81	$3d_5-4Y$.81	1	10,420.52	9,593.83	$2p_7-4d_5$.74
25	9,211.38	10,853.16	$3d_5-4X$.14	8	10,484.83	9,534.98	$2p_7-4d_6$.98
1	9,216.51	10,847.12	$3d_5-6Y$.18	6	10,507.91	9,514.03	$3d_5-5V$.02
5	9,222.39	10,840.20	$2p_8-4d_1''$.19	10	10,515.15	9,507.43	$3d_5-2p_2$.44
3	9,245.18	10,813.48	$2p_8-3s_4$.465	40	10,527.84	9,496.03	$2p_6-4d_3$	5.99
30	9,301.95	10,747.49	$3d_1'-4W$.40	4	10,549.75	9,476.30	$3d_2-5Y$.30
40	9,306.64	10,742.07	$1s_3-3p_{10}$.092	25	10,706.77	9,337.33	$2p_6-4d_5$.31
3	9,334.08	10,710.49	$3d_1'-4U$.45	15	10,758.85	9,292.13	$3d_5-2p_1$.07
10	9,374.02	10,664.86	$3d_1'-4Z$.90	100	10,838.37	9,223.95	$1s_1-2p_{10}$.986
100	9,374.76	10,664.02	$3d_1'-4T$.02	10	10,895.39	9,175.68	$3d_5-2p_3$.72
60	9,412.01	10,621.81	$2p_8-4d_4$.83	10	11,085.39	9,018.42	$3d_1''-4W$.44
20	9,441.46	10,588.68	$2p_9-4d_3$.66	2	11,127.29	8,984.46	$3d_1''-4V$.46
20	9,442.68	10,587.31	$3d_5-4V$.30	1	11,140.90	8,973.48		
80	9,445.34	10,584.33	$3d_5-4U$.33					
4	9,487.76	10,537.01							
40	9,497.07	10,526.68	$3d_5-4Y$.69					
10	9,505.78	10,517.03	$3d_5-4X$.02					
200	9,513.379	10,508.633	$2p_8-4d_1'$.633					
20	9,535.14	10,429.96	$2p_9-4d_5$.98					

IV. DISCUSSION OF Xe I TERMS

The numerical values of previously known terms have been subject to only slight revisions. The new term table incorporates the changes due to more precise data, the extensions of series, new low terms and changes in the assignment of quantum numbers. The values of the $1s$, $2p$, and $3p$ terms as determined from interference measurements (RP245) form the basis of the present table. The numerical values of a considerable number of terms can now be completely determined from interference measurements. These are given to three places of decimals. Table 2 gives the final set of terms and effective quantum numbers. In calculating the effective quantum numbers for non-Ritzian terms, the value of the displacement constant A was taken as 10,540, representing the level separation of the ground doublet ($^2P_{1/2}, ^2P_{3/2}$) identified⁹ in the Xe II spectrum.

⁹ C. J. Humphreys, T. L. deBruin, and W. F. Meggers, B. S. Jour. Research, vol. 6 (RP275), p. 237, 1931.

Elec- troh	9		10		11		12	
<i>p</i>								
<i>s</i>	1, 102. 57	9. 9764	909. 3	10. 98				
<i>p</i>	1, 227. 12	9. 4565	1, 003. 75	10. 4559	835. 2	11. 46		
	1, 216. 99	9. 4958	996. 75	10. 4926				
	1, 199. 80	9. 5112						
<i>d</i>	1, 512. 34	8. 5183	1, 211. 09	9. 5190	991. 76	10. 5191	826. 99	11. 5193
	1, 499. 413	8. 5549	1, 202. 24	9. 5539	985. 30	10. 5534		
	1, 505. 27	8. 5383	1, 206. 81	9. 5319				
	1, 518. 73	8. 5003	1, 217. 67	9. 4932				
	1, 410. 12	8. 8216	1, 139. 47	9. 8135				
	1, 485. 86	8. 5939	1, 192. 55	9. 5926	978. 32	10. 5909		
	1, 475. 33	8. 6245	1, 184. 83	9. 6238	972. 39	10. 6232		
<i>f</i>	1, 371. 26	8. 9457						
	1, 370. 62	8. 9478						
	1, 368. 12	8. 9560	1, 107. 27	9. 9552	914. 64	10. 9535		
	1, 368. 26	8. 9556	1, 107. 42	9. 9545	914. 67	10. 9533		
	1, 364. 16	8. 9690	1, 104. 40	9. 9681	912. 38	10. 9671		
	1, 363. 50	8. 9712	1, 102. 60	9. 9762				
	1, 360. 71	8. 9804						

TABLE 3.—*Determination of absolute term values of XeI*

$$2p_3-md'_4 \quad (5p^3D_3-5d^3F_4)$$

$$\nu = A - \frac{R}{[m + \mu + \alpha(-1 - \nu)]^2}$$

$$R_{\infty} = 109,736.965$$

$$A = 19,430.838$$

$$\mu = +0.5208392314$$

$$\alpha = -0.000001534768614$$

<i>m</i>	Int.	$\lambda_{\text{obs.}}$	$\nu_{\text{obs.}}$	$\nu_{\text{calc.}}$	$\nu_{\text{calc.}} - \nu_{\text{obs.}}$
2	-----	-----	[1,793.60]	1,784.900	+8.70
3	200	9,513.32	10,508.70	10,509.189	-.48
4	500	7,119.598	14,041.872	14,041.873	-.001
5	300	6,318.062	15,823.272	15,823.272	-.000
6	100	5,934.170	16,846.900	16,846.946	-.046
7	70	5,716.255	17,489.131	17,489.217	-.086
8	40	5,579.276	17,918.51	17,918.584	-.07
9	6h	5,437.026	18,219.76	18,219.760	-.000
10	2h	5,421.756	18,439.10	18,439.144	-.04
11	1h-	5,373.74	18,603.85	18,603.886	-.04

$$2p_3 = 19,430.838$$

$$1s_5 - 2p_3 = 11,335.515$$

$$\therefore 1s_5 = 30,766.353$$

The absolute value of the term $1s_5$ has been redetermined using the more accurate wave length data now available and a new value of the Rydberg constant for xenon calculated from the value of R_{∞} given by Birge.¹⁰ The new value of $1s_5$ differs very slightly from the old, 30,766.353 as compared with 30,766.98 previously calculated. This revised term value represents the precision attainable with the data selected for the determination. It is still subject to a variation of from 1 to 2 wave number units, depending on the series chosen for computation. As in the work previously reported, we have estimated $1s_5$ from the limit of the series $2p_3-md'_4$ and the combination $1s_5-2p_3$. All other terms are fixed relative to $1s_5$, so that the error in absolute value is the same for all terms.

The sequence ms_5 remains unchanged except that it has been extended to the tenth member. Attempts to find $2s_5$ from combinations in the infra-red region have not been successful, and it seems probable that the required combinations are beyond our range of photographic observations. The sequence ms_4 has also been extended. The term 13,943.93 previously given as $2s_4$ is now interpreted as $3d_2$. The reason for the new assignment will be discussed in connection with the d -series. The old terms, $2s_3$ and $2s_2$, originally obtained from the interpretation of the ultra-violet data of Abbink and Dorgelo¹¹ are not confirmed by expected combinations in the range of the present observations and are not retained. The term, 4,215.65, is a new term, interpreted as $2s_2$ by Rasmussen, and is confirmed by observed combinations.

No changes have been made in the interpretation of the p -series. Higher members have been found in the case of all sequences which form Ritz series. The sequence mp_3 is the longest; it has been extended to $11p_3$.

A considerable number of changes have been introduced in the interpretation of the d -series. Several new terms have been found

¹⁰ R. T. Birge, Reviews of Modern Physics, vol. 1, p. 1, 1929.

¹¹ J. H. Abbink and H. B. Dorgelo, Zeit. f. Phys., vol. 47, p. 221, 1928.

and all Ritzian series except md_6 have been extended. In the work previously reported, the failure to find all expected combinations, due to the faintness of the sources used and the lack of the superior sensitizers now available for the infra-red region, left the assignment of inner quantum numbers uncertain in many cases. With the more extended and accurate data now available these ambiguities have been eliminated and it is believed that the j -values are now determined correctly for all the d -terms. In the work of Rasmussen referred to above, new interpretations have been placed on many of the old terms and several additional d -terms have been found. The present work entirely confirms Rasmussen's assignment of j -values, and we agree also on interpretation of the terms in most cases.

The md_6 -sequence has been altered only by the substitution of the new term 9,342.88 for $4d_6$ in place of 9,125.44, which by virtue of its j -value 2 is assigned to $4d_3$.

The md_4' sequence is left unchanged. It is the longest as well as the most regular of the d -series, and has been extended in the present work to $12d_4'$. Attention is again called to Table 3 in which a Ritz formula is applied to the series $2p_8-md_4'$. The only large departure from the calculated position of a line is in the case of the photographically inaccessible infra-red line $2p_8-3d_4'$, at $1,793.60\text{ cm}^{-1}$. The position of this line is predicted accurately from the known values of both terms, and differs from the position given by the series formula by 8.70 cm^{-1} .

The first term only of the md_4 sequence has been changed. The term 16,863.42, previously assigned to $3d''$, is now designated $3d_4$; 17,511.12, formerly called $3d_4$, becomes the first term of the new sequence md_3 with inner quantum number 2. Subsequent terms of md_3 are assembled from members of other d -series incorrectly interpreted. The second number $4d_3$ is 9,125.44, formerly $4d_6$, the third $5d_3$ is 5,112.357, formerly $5d''$, whereas the remaining terms of md_3 are taken from the old md_5 sequence.

The md_5 sequence begins as before with 17,847.24, as $3d_5$. The term 9,284.12, formerly $4d_2$, is now assigned to $4d_5$. The term 5,705.605 was formerly designated as $3s_1'''$. It is now certain that its j -value is 1. Rasmussen designated it as $5d_5$. We give it the same classification with reservations, since its interpretation as the missing non-Ritzian term $3s_1'$ seems equally probable. The present classification introduces a sharp downward inflection in the graph of the effective quantum numbers, which is not paralleled for Kr I. In any case either $5d_5$ or $3s_1'$ has not been found.

With the exception of the first and third members, the md_2 sequence is made up of newly found terms. The term 13,943.93 is now designated as $3d_2$ instead of $2s_4$. This assignment is supported by the effective quantum number which fits better into this d -series than into the s_4 -series where an unexpected deviation from an otherwise fairly regular series would occur. Furthermore, if $2s_4$ is found, $2s_5$ should also be observed, since the combinations would be displaced to shorter wave lengths than those of $2s_4$. The combination of 13,944 with the ground term giving the line, $83,891.2\text{ cm}^{-1}$, found in Abbink and Dorgelo's data is permitted with either interpretation. The term 5,119.845, designated $5d_2$, was formerly $5d_5$. Three new terms, 7,801.75, 3,148.46, 2,335.41, have been interpreted by Rasmussen as

$4d_2$, $6d_2$, and $7d_2$, respectively. We have now extended the series to $10d_2$.

The first three members of the d_1'' sequence have been replaced. The first member 15,908.28, now $3d_1''$, was formerly $3d_1'$. The term 8,289.42, appearing in the original table, is not supported by observed combinations and is obviously spurious. It is replaced by 8,590.65, designated $4d_1''$ instead of $4d_5$ as before. The third term, 5,155.368, has been changed to $5d_1''$ from $5d_2$. The original interpretation has been retained for subsequent members of the series.

A new term 15,403.68 has been interpreted by Rasmussen as $3d_1'$. The md_1' sequence is otherwise unchanged.

Of the original group of d -type non-Ritzian terms, two are retained; 6,087 is designated $3s_1'''$ instead of $3s_1''''$, and 6,386, now $3s_1''''$, was formerly $3s_1''$. These changes were suggested by Rasmussen to make the notation uniform among the rare gas spectra. A new term 6,681 becomes $3s_1''$. Our old $3s_1'$ term is not confirmed. The only remaining possibility for $3s_1'$ is 5,705 as explained above.

Changes in the f -type or hydrogenlike terms have been limited to small corrections and extensions of series, except that one entirely new sequence mT has been found. A discrepancy between observed and calculated wave numbers, too great to be accounted for by errors of observation, was found in the case of the series of lines formerly designated $3d_4'-mZ$. Examination of the first and second members of the series at 9,374 and 7,584 Å with the Rowland grating revealed a companion line in each case in the calculated position of the corresponding member of the mZ series. The stronger series was, therefore, given a separate designation mT and the weaker retained as mZ . The separation is 0.9 cm^{-1} for the first member 0.7 cm^{-1} for the second, the two series merging in the higher members. Members beyond the second can not be resolved not only because the wave number separation is smaller, but also since the wave length difference diminishes much more rapidly as one goes to shorter wave lengths.

The data here presented, because of greater completeness, accuracy, and the elimination of errors, should be regarded as superseding those given in RP115.

WASHINGTON, November 3, 1932.



BUREAU OF STANDARDS JOURNAL OF RESEARCH

CONTENTS OF RECENT NUMBERS—Continued

October, 1932 (Vol. 9, No. 4)

- RP482. The synthesis, purification, and certain physical constants of the normal hydrocarbons from pentane to dodecane, of *n*-amyl bromide and of *n*-nonyl bromide. B. J. Mair. Price 5 cents.
- RP483. A study of some ceramic bodies of low absorption maturing at temperatures below 1,000° C. R. F. Geller and D. N. Evans. Price 5 cents.
- RP484. The determination of magnesia in phosphate rock. James I. Hoffman. Price 5 cents.
- RP485. Collisions of the first and second kind in the positive column of caesium discharge. F. L. Mohler. Price 5 cents.
- RP486. The areas and tensile properties of deformed concrete-reinforcement bars. A. H. Stang, L. R. Sweetman, and C. Gough. Price 5 cents.
- RP487. A calorimetric method for determining the intrinsic energy of a gas as a function of the pressure. Edward W. Wasburn. Price 5 cents.
- RP488. The photographic emulsion: Variables in sensitization by dyes. Burt H. Carroll and Donald Hubbard. Price 5 cents.
- RP489. A method for the separation of rhodium from iridium and the gravimetric determination of these metals. Raleigh Gilchrist. Price 5 cents.
- RP490. The isoelectric point of silk. Milton Harris. Price 5 cents.
- RP491. Effective applied voltage as an indicator of the radiation emitted by an X-ray tube. Lauriston S. Taylor, George Singer, and C. F. Stoneburner. Price 5 cents.
- RP492. Column curves and stress-strain diagrams. William R. Osgood. Price 5 cents.

November, 1932 (Vol. 9, No. 5)

- RP493. The variation with angle of emission of the radiation from metals bombarded by slow electrons. C. Boeckner.
- RP494. Nitrogen content of some standard sample steels. J. G. Thompson and E. H. Hamilton.
- RP495. Analytical methods for the determination of levulose in crude products. R. F. Jackson, J. A. Mathews, and W. D. Chase.
- RP496. Determination of alumina and silica in steel by the hydrochloric acid residue method. J. G. Thompson and J. S. Acken.
- RP497. An analysis of lanthanum spectra (La I, La II, La III). Henry Norris Russell and William F. Meggers.
- RP498. Analysis of weighted silk. Ralph T. Mease.
- RP499. The heat of formation of hydrogen chloride and some related thermo-dynamic data. Frederick D. Rossini.
- RP500. Thermal expansion of lead. Peter Hidnert and W. T. Sweeney.

December, 1932 (Vol. 9, No. 6)

- RP501. Isolation of the three xylenes from an Oklahoma petroleum. Joseph D. White and F. W. Rose, jr.
- RP502. Characteristic equations of vacuum and gas-filled tungsten-filament lamps. L. E. Barbrow and J. Franklin Meyer.
- RP503. The calorimetric determination of the intrinsic energy of gases as a function of the pressure. Data on oxygen and its mixtures with carbon dioxide to 40 atmospheres, at 28° C. Frederick D. Rossini and Mikkel Frandsen.
- RP504. Shear tests of reinforced brick masonry beams. D. E. Parsons, A. H. Stang, and J. W. McBurney.
- RP505. The measurement of low-voltage X-ray intensities. Lauriston S. Taylor and C. F. Stoneburner.
- RP506. A multi-range potentiometer and its application to the measurement of small temperature differences. H. B. Brooks and A. W. Spinks.
- RP507. The density of some soda-lime-silica glasses as a function of the composition. F. W. Glaze, J. C. Young, and A. N. Finn.
- RP508. Coiled-filament resistance thermometers. C. H. Meyers.
- RP509. Deflection of cosmic rays by a magnetic field. L. F. Curtiss.
- RP510. The system: $\text{CaO-P}_2\text{O}_5$. Elmer T. Carlson.
- RP511. A metal-connected glass electrode. M. R. Thompson.

Research papers, where price is given, are already available, and others will be available in the near future as separates for purchase from the Superintendent of Documents, United States Government Printing Office, Washington, D. C.



